NSI at Neutrino Factories: correlations & degeneracies

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Based on a work in collaboration with: A. Donini, J. López-Pavón and H. Minakata

 $\mathcal{L}_{eff} = \mathcal{L}^{SM} + \mathcal{L}_{\nu}^{mass} + \sum c_i O_i^{p,d,f}$

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NSI@production

 $O^{p}: (\epsilon^{p}_{e\alpha}(\bar{\mu}\gamma^{\mu}_{L}\nu_{\mu})(\bar{\nu}_{\alpha}\gamma_{\mu L}e) \qquad \mu^{-} \to e^{-}\nu_{\mu}\bar{\nu}_{\alpha}$

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NSI@detection

$$O^d: \epsilon^d_{\mu\alpha} (\bar{\nu}_{\alpha} \gamma^{\mu}_L \mu) (\bar{d} \gamma_{\mu L} u)$$

$$\left(\
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NSI@propagation

$$O^f: \epsilon^f_{\alpha\beta}(\bar{\nu}_{\alpha}\gamma^{\mu}_L\nu_{\beta})(\bar{f}\gamma_{\mu}f)$$

$$u_lpha f o
u_eta f$$

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 $\mathcal{L}_{eff} = \mathcal{L}^{SM} + \mathcal{L}_{\nu}^{mass} + \sum c_i O_i^{p,d,f}$

Near Detectors:

S. Antusch *et al*, arXiv:1005.0756 [hep-ph] MINSIS workshop report, arXiv:1009.0476 [hep-ph] NSI@production

$$\mu^- \to e^- \nu_\mu \bar{\nu}_\alpha$$

NSI@detection

$$\nu_{\alpha}N \to \mu^{-}N'$$

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$$\mu^- o e^-
u_\mu ar
u_lpha$$

NSI@detection

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ightarrow \mu^- N'$$

NSI@propagation

$$u_{lpha} f
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Effective matter potential: Far detectors

 $A^{NSI} = A \begin{pmatrix} 1 + \epsilon_{ee}^m & \epsilon_{e\mu}^m & \epsilon_{e\tau}^m \\ \epsilon_{e\mu}^{m*} & \epsilon_{\mu\mu}^m & \epsilon_{\mu\tau}^m \\ \epsilon_{e\tau}^{m*} & \epsilon_{\mu\tau}^{m*} & \epsilon_{\tau\tau}^m \end{pmatrix}$

T. Kikuchi, H. Minakata, S. Uchinami arXiv:0809.3312v2 [hep-ph]



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Brief review to analytical dependences:

Diagonal
$$\begin{cases} P_{\alpha\beta}(\epsilon_{ee} - \epsilon_{\tau\tau}) \sim \mathcal{O}(\varepsilon^3) !! \\ P_{\alpha\beta}(\epsilon_{\mu\mu} - \epsilon_{\tau\tau}) \sim \mathcal{O}(\varepsilon^2) \end{cases}$$

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$$A^{NSI} = A \begin{pmatrix} 1 + \epsilon^m_{ee} & \epsilon^m_{e\mu} & \epsilon^m_{e\tau} \\ \epsilon^{m*}_{e\mu} & \epsilon^m_{\mu\mu} & \epsilon^m_{\mu\tau} \\ \epsilon^{m*}_{e\tau} & \epsilon^{m*}_{\mu\tau} & \epsilon^m_{\tau\tau} \end{pmatrix}$$

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$$\begin{array}{l} \text{Off-} \\ \text{diagonal} \\ \text{sector} \end{array} \left\{ \begin{array}{l} P_{e\mu,e\tau} = P_{e\mu,e\tau}^{std} + \mathcal{O}(\varepsilon^2) & (\epsilon_{e\mu}, \ \epsilon_{e\tau}, \ \epsilon_{\mu\tau}) \\ P_{\mu\mu,\mu\tau} = P_{\mu\mu,\mu\tau}^{std} + \mathcal{O}(\epsilon_{\mu\tau}) + \mathcal{O}(\varepsilon^2) \end{array} \right.$$

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- WARNINGS!
 - Suppression with scale of New Physics
 - Many parameters to be introduced at once

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- WARNINGS!
 - Suppression with scale of New Physics
 - Many parameters to be introduced at once
- Up to now, no correlations studied in literature
 - MonteCUBES allows to introduce all parameters at once (M.Blennow, E. Fernández-Martínez; arXiv:0903.3985 [hep-ph])

 Astonishing sensitivities to standard oscillation parameters

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- Long baseline
- High energies



Large matter effects!

Multi-channel facility

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- But...what if θ_{13} is measured soon?

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Large matter effects!

- Multi-channel facility
- But...what if θ_{13} is measured soon?
 - Open possibility: re-optimization of NF to search for New Physics?

Setups

• IDS25:

- 25 GeV muons;
- Two 50 kton MIND detectors (arXiv:1004.0358 [hep-ex]):
 - @4000 km: good for CP
 - @7500 km: good for $heta_{13}$ and hierarchy (MB)
- 5×10^{20} useful muon decays/year/baseline/polarity

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IDS50: 50 GeV upgrade of the IDS25

Setups

But the NF is multi-channel!
So we will study a 3rd setup:
1B50:

- 50 GeV muons:
- A composite detector @ 4000 km:
 - 50 kton MIND to detect muons;
 - 4 kton MECC to detect taus (arXiv:hep-ph/0305185).
- Double flux: 10²¹ useful muon decays/year/polarity

1^{s} question: Does the sensitivity to θ_{13} change in presence of NSI?

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YES

Sensitivity to θ_{13} in <u>presence</u> of NSI



No correlation at all with $\epsilon_{\mu\tau}$ Worsening exclusively due to $\epsilon_{\alpha\alpha}$

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Strong correlation due to simultaneous appearance in golden channel

Sensitivity to θ_{13} in <u>presence</u> of NSI



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2rd question: Will we actually be able to see NSI?

Blennow, Meloni, Ohlsson, Terranova, Westerberg, arXiv:0804.2744 [hep-ph]

Kopp, Ota, Winter, arXiv:0804.2261 [hep-ph]





(Marginalization performed over all standard parameters)



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Sensitivity to $\epsilon_{\mu \tau}$



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Sensitivity to $\epsilon_{\mu \tau}$

• Good news: No correlation with θ_{13} , $\epsilon_{\alpha\beta}$ Mild correlation with $\epsilon_{\alpha\alpha}$

 Same result for all setups under study.

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Sensitivity to $\epsilon_{e\mu}, \epsilon_{e\tau}$



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3rd question: What about CP violation?



99 % CL 2 d.o.f. (No marginalization performed, but over $\theta_{13}, \, \delta$)







Conclusions

- Generically, we conclude that higher energy setups are better to study NSI;
- How do NSI affect θ_{13} sensitivity?
 - No correlations with $\epsilon_{\mu\tau}$ are observed;
 - Mild correlations with $\epsilon_{\alpha\alpha}$;
 - Strong correlations with $\epsilon_{e\mu}, \epsilon_{e\tau}$.
- Diagonal NSI parameters:
 - Sizable effects due to $\theta_{13} \neq 0$; $\delta \theta_{23} \neq 0$,
 - Sensitivity $\mathcal{O}(10^{-1})$ for $(\epsilon_{ee} \epsilon_{\tau\tau})$; $\mathcal{O}(10^{-2})$ for $(\epsilon_{\mu\mu} - \epsilon_{\tau\tau})$.

Conclusions

- Off-diagonal parameters:
 - $\epsilon_{e\mu}$: higher energies are the key
 - $\epsilon_{e\tau}$: the MB is the key factor

Sensitivities of $\mathcal{O}(10^{-3})$ are achievable

- $\epsilon_{\mu\tau}$: independent of setup. Linear dependence on real part gives rise to sensitivities ranging from $10^{-3} 10^{-2}$
- CP violation:
 - CP violation exclusively due to NSI could be measured for vanishing θ_{13} in a 35% of the phase space for $|\epsilon_{e\mu}| = 0.001; |\epsilon_{e\tau}| = 0.01.$