# Resolving the Octant $\theta_{23}$ Degeneracy by Neutrino Oscillation Experiments 

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## Based on the works

K. Hiraide, H. Minakata, T. Nakaya, H. Sugiyama, W.J.C.Teves, R. Zukanovich Funchal, HN, hep-ph/060I258 [PRD73, 093008 (2006)]
T. Kajita, H. Minakata, S. Nakayama, HN, to appear
@NOW2006, September 15, 2006

## Outline

- Introduction
- Resolving of the $\theta_{23}$ Octant Degeneracy by (I) accelerator (superbeam) and reactor
- Resolving of the $\theta_{23}$ Octant Degeneracy by
(2) accelerator (superbeam) with 2 identical detectors
- Summary


## Introduction

$\theta_{23}$ octant degeneracy is a part of the so called 8 fold parameter degeneracy we encounter when we try to determine neutrino mixing paramters by oscillation experiments
(I) $\theta_{23}$ octant degeneracy (Fogli and Lisi, I996)
(2) Intrinsic $\theta_{13} \delta$ degeneracy (Burguet-Castell et al, 200I)
(3) sign of $\Delta \mathrm{m}_{13}^{2}$ degeneracy (Minakata, $\mathrm{HN}, 200 \mathrm{I}$ )


8 fold degeneracy (Barger et al, 2002)

## Currently, we don't know if $\theta_{23}$ is maximal

From SK atmospheric neutrino data

$$
\begin{aligned}
\sin ^{2} 2 \theta_{23} & >0.92 @ 90 \% \mathrm{CL} \\
0.34 & <\sin ^{2} \theta_{23}<0.64
\end{aligned}
$$

From Theory Consideration: $\mu-\tau$ symmetry

> Typically predict $\theta_{23}=\pi / 4$ and $\theta_{13}=0$
> in the symmetry limit

## What is the $\theta_{23}$ Octant degeneracy?

## Fogli and Lisi, PRD 54, 3667 (1996)

Let us first consider the disapperance mode $\nu_{\mu} \rightarrow \nu_{\mu}$
One $\Delta \mathrm{m}^{2}$ dominance approximation and vanishing $\theta_{13}$

$$
\mathrm{P}\left(v_{\mu} \rightarrow v_{\mu}\right) \approx 1-\sin ^{2} 2 \theta_{23} \sin ^{2}\left(\frac{\Delta \mathrm{~m}^{2}}{4 \mathrm{E}} \mathrm{~L}\right)
$$

Disappearance mode can determine $\sin ^{2} 2 \theta_{23}$

$$
\sin ^{2} \theta_{23}=\frac{1}{2}\left[1 \pm \sqrt{1-\sin ^{2} 2 \theta_{23}}\right]
$$

For example, $\sin ^{2} 2 \theta_{23}=0.96 \rightarrow \sin ^{2} \theta_{23}=0.4$ or 0.6

## To be more precise...

$$
\mathrm{P}\left(v_{\mu} \rightarrow v_{\mu}\right) \equiv 1-\sin ^{2} 2 \theta_{\mathrm{eff}} \sin ^{2}\left(\frac{\Delta \mathrm{~m}^{2} \mathrm{eff}}{4 \mathrm{E}} \mathrm{~L}\right)
$$

$$
\sin ^{2} 2 \theta_{\text {eff }} \approx 4\left|\mathrm{U}_{\mu 3}\right|^{2}\left(1-\left|\mathrm{U}_{\mu 3}\right|^{2}\right) \begin{aligned}
& \mathrm{T} 2 \mathrm{~K} I \text { can measure this quantity } \\
& \text { with } \approx 1 \% \text { accuracy }
\end{aligned}
$$

For a given value of $\sin ^{2} 2 \theta_{\text {eff }}$

$$
\sin ^{2} \theta_{23} \approx \frac{1}{2}\left[1 \pm\left(1+\sin ^{2} \theta_{13}\right) \sqrt{1-\sin ^{2} 2 \theta_{\text {eff }}}\right]
$$

$$
\sin ^{2} 2 \theta_{\text {eff }}=0.96 \rightarrow \sin ^{2} \theta_{23} \approx(0.4 \text { or } 0.6)+\underline{0.1 \sin ^{2} \theta_{13}}
$$

## What will happen if we add $v_{\mu} \rightarrow v_{e}$ ?

One $\Delta m^{2}$ dominance approximation

$$
\mathrm{P}\left(v_{\mu} \rightarrow v_{\mathrm{e}}\right) \approx \sin ^{2} \theta_{23} \sin ^{2} 2 \theta_{13} \sin ^{2}\left(\frac{\Delta \mathrm{~m}^{2}}{4 \mathrm{E}} \mathrm{~L}\right)
$$

Appearance mode can determine $\sin ^{2} \theta_{23} \sin ^{2} 2 \theta_{13}$

$$
\left(\text { Not } \theta_{23} \text { and } \theta_{13}\right. \text { separately !) }
$$

Unless we know $\theta_{13}$ determined by some OTHER experiment, we can not distinguish 2 value of $\sin ^{2} \theta_{23}$ obtained by the disappearance mode!

Two solutions of $\left(\theta_{23}, \theta_{13}\right)$

$$
\sin ^{2} \theta_{23}^{\text {fake }} \sin ^{2} 2 \theta_{13}^{\text {fake }} \approx \sin ^{2} \theta_{23}^{\text {true }} \sin ^{2} 2 \theta_{13}^{\text {true }}
$$

## Some Examples of Degenerate Solutions



Assume T2K Phase II: 4MW (2yrv + 6 yr $\bar{v})$, HK@Kamioka

## How Can We Resolve This Degeneracy?

## First Possible Strategy

## Combine Reactor Data

 Minakata et al, PRD 68, 033017 (2003) See also Fogli and Lisi, PRD 54, 3667 (1996)$\mathrm{P}\left(\nu_{\mathrm{e}} \rightarrow \nu_{\mathrm{e}}\right) \approx 1-\sin ^{2} 2 \theta_{13} \sin ^{2}\left(\frac{\Delta \mathrm{~m}_{13}^{2}}{4 \mathrm{E}} \mathrm{L}\right)+\mathrm{O}\left(\sin ^{2} 2 \theta_{13} \Delta \mathrm{~m}_{12}^{2}\right)$
Reactor with $L \approx 1-1.5 \mathrm{~km}$ can provide a clean measurement of $\theta_{13}$ free from degeneracy

Since 2 degenerate solutions correspond to different values of $\theta_{13}$, it will be possible to eliminate one of the solutions

## Experimental Setup and Assumptions

(I) T2K Phase II: 4 MW 2.5 deg. OA beam ( $2 \mathrm{yr} v+6 \mathrm{yr} \overline{\mathrm{v}}$ ) HK 0.54 Mt , appearance and disappearance modes
(2) High Statistics/Sensitivity Reactor Experiment with an exposure of $\sim 10 \mathrm{GW} \cdot \mathrm{kt} \cdot \mathrm{yr}, \mathrm{L}=1.5 \mathrm{~km}$ beyond Double CHOOZ
Inputs
$\Delta \mathrm{m}_{23}^{2}= \pm 2.5 \times 10^{-3} \mathrm{eV}^{2}$
$\Delta \mathrm{m}^{2}{ }_{12}=8 \times 10^{-5} \mathrm{eV}^{2}$
$\sin ^{2} \theta_{12}=0.31$
$\delta=0$

## $\chi^{2}$ definitions

(1) $v_{\mu} \rightarrow v_{e}$ appearance mode

$$
\chi_{\mathrm{app}}^{2} \equiv \frac{\left(N_{\mathrm{sig}}^{\mathrm{obs}}+N_{\mathrm{BG}}^{\mathrm{obs}}-N_{\mathrm{sig}}^{\mathrm{theo}}-N_{\mathrm{BG}}^{\mathrm{theo}}\right)^{2}}{N_{\mathrm{sig}}^{\mathrm{obs}}+N_{\mathrm{BG}}^{\mathrm{obs}}+\left(\sigma_{\mathrm{sig}} N_{\mathrm{sig}}^{\mathrm{obs}}\right)^{2}+\left(\sigma_{\mathrm{BG}} N_{\mathrm{BG}}^{\mathrm{obs}}\right)^{2}}
$$

(2) $v_{\mu} \rightarrow v_{\mu}$ disappearance mode
$\chi_{\mathrm{dis}}^{2} \equiv \min _{\alpha_{\mathrm{sig}, \alpha_{\mathrm{BG}}}} \sum_{i} \frac{\left[N_{i}^{\mathrm{obs}}+N_{i, \mathrm{BG}}^{\mathrm{obs}}-\left(1+\alpha_{\mathrm{sig}}\right) N_{i}^{\mathrm{theo}}-\left(1+\alpha_{\mathrm{BG}}\right) N_{i, \mathrm{BG}}^{\mathrm{theo}}\right]^{2}}{N_{i}^{\mathrm{obs}}+N_{i, \mathrm{BG}}^{\mathrm{obs}}}+\left(\frac{\alpha_{\mathrm{sig}}}{\sigma_{\mathrm{sig}}}\right)^{2}+\left(\frac{\alpha_{\mathrm{BG}}}{\sigma_{\mathrm{BG}}}\right)^{2}$,
(3) $\bar{v}_{\mathrm{e}} \rightarrow \bar{v}_{\mathrm{e}}$ reactor disappearance mode

$$
\sigma_{\mathrm{sig}}=\sigma_{\mathrm{BG}}=2 \%: \text { sytematic erros }
$$

$\chi_{\text {reac }}^{2} \equiv \min _{\alpha^{\prime} \mathrm{s}} \sum_{a=f, n}\left[\sum_{i=1}^{17}\left\{\frac{\left(N_{a i}^{\mathrm{theo}}-\left(1+\alpha_{i}+\alpha_{a}+\alpha\right) N_{a i}^{\mathrm{obs}}\right)^{2}}{N_{a i}^{\mathrm{obs}}+\sigma_{\mathrm{db}}^{2}\left(N_{a i}^{\mathrm{obs}}\right)^{2}}+\frac{\alpha_{i}^{2}}{\sigma_{\mathrm{Db}}^{2}}\right\}+\frac{\alpha_{a}^{2}}{\sigma_{\mathrm{dB}}^{2}}\right]+\frac{\alpha^{2}}{\sigma_{\mathrm{DB}}^{2}}$,
(a) $\sigma_{D B}=\sigma_{D b}=2 \% \quad \sigma_{d B}=\sigma_{d b}=0.5 \%$ : Conservative choice of sys. error
(b) $\sigma_{\mathrm{DB}}=\sigma_{\mathrm{Db}}=1 \% \quad \sigma_{\mathrm{dB}}=\sigma_{\mathrm{db}}=0.2 \%$ : Optimistic choice of sys. error

D(B): Correlated between Detectors(Bins) d(b): Uncorrelated between Detectors(Bins)

## Impact of adding Reactor data



## Impact of adding Reactor data



## Expected Sensitivity: Regions of paramters where the hierarchy can be determined



## Expected Sensitivity: Regions of paramters where the hierarchy can be determined



## Second Possible Strategy

T. Kajita, H. Minakata, S. Nakayama and HN, to appear

## Superbeam with 2 detector system

M. Ishitsuka, T. Kajita, H. Minakata, HN, hep-ph/0504026

$$
\begin{aligned}
& P\left[\nu_{\mu}\left(\bar{\nu}_{\mu}\right) \rightarrow \nu_{\mathrm{e}}\left(\bar{\nu}_{e}\right)\right]=\left.c_{23}^{2}\right) \sin ^{2} 2 \theta_{12}\left(\frac{\Delta m_{21}^{2} L}{4 E}\right)^{2} \longleftarrow \text { Solar Term }=\mathrm{P}_{\text {solar }} \\
&+\sin ^{2} 2 \theta_{13} s_{23}^{2}\left[\sin ^{2}\left(\frac{\Delta m_{31}^{2} L}{4 E}\right)-\frac{1}{2} s_{12}^{2}\left(\frac{\Delta m_{21}^{2} L}{2 E}\right) \sin \left(\frac{\Delta m_{31}^{2} L}{2 E}\right)\right. \\
&\left. \pm\left(\frac{4 E a}{\Delta m_{31}^{2}}\right) \sin ^{2}\left(\frac{\Delta m_{31}^{2} L}{4 E}\right) \mp \frac{a L}{2} \sin \left(\frac{\Delta m_{31}^{2} L}{2 E}\right)\right] \\
&+2 J_{r}\left(\frac{\Delta m_{21}^{2} L}{2 E}\right)\left[\cos \delta \sin \left(\frac{\Delta m_{31}^{2} L}{2 E}\right) \mp 2 \sin \delta \sin ^{2}\left(\frac{\Delta m_{31}^{2} L}{4 E}\right)\right] \\
& \equiv \mathrm{P}_{\text {solar }}+\mathrm{P}_{\mathrm{atm}}
\end{aligned}
$$

Magnitude of the Solar term: $P_{\text {solar }} \equiv \cos ^{2} \theta_{23} \sin ^{2} 2 \theta_{12} \sin ^{2}\left(\frac{\Delta \mathrm{~m}_{12}^{2}}{4 \mathrm{E}} \mathrm{L}\right)$


Effect of solar term is very small at Kamioka ( $\mathrm{L}=295 \mathrm{~km}$ ) but sizable at Korea ( $\mathrm{L}=1050 \mathrm{~km}$ )

## Experimental Setup and Assumptions

T2KK (Tokai to Kamioka-Korea)

0.27 Mt detector @Kamioka (L=295 km) 2 Identitical detectors at the same 0.27 Mt detector @Korea (L=I050km) off axis angle
Appearance and disappearance modes

$$
\Delta \mathrm{m}_{23}^{2}= \pm 2.5 \times 10^{-3} \mathrm{eV}^{2} \quad \Delta \mathrm{~m}_{12}^{2}=8 \times 10^{-5} \mathrm{eV}^{2} \quad \sin ^{2} \theta_{12}=0.3 \mathrm{I}
$$

## Impact of the Solar Term



## $\chi^{2}$ definition

$\binom{v}{\bar{v}} \times\binom{$ kamioka }{ korea }$=4$ combinations

$$
\begin{gathered}
v_{\left.\mathrm{e}^{\left(\bar{v}_{\mathrm{e}}\right.}\right) \text { event (5 bins) }}^{\sigma_{i}^{2}}\left(\sum_{i=1}^{5} \frac{\left(N(e)_{i}^{\mathrm{obs}}-N(e)_{i}^{\exp }\right)^{2}}{\left.\sigma_{\mu}\right) \text { event (20 bins) }}+\sum_{i=1}^{20} \frac{\left(N(\mu)_{i}^{\mathrm{obs}}-N(\mu)_{i}^{\exp }\right)^{2}}{\sigma_{i}^{2}}\right)+\sum_{j=1}^{7}\left(\frac{\epsilon_{j}}{\tilde{\sigma}_{j}}\right)^{2} \\
N(e)_{i}^{\exp }=N_{i}^{\mathrm{BG}} \cdot\left(1+\sum_{j=1,2,7} f(e)_{j}^{i} \cdot \epsilon_{j}\right)+N_{i}^{\mathrm{signal}} \cdot\left(1+\sum_{j=3,7} f(e)_{j}^{i} \cdot \epsilon_{j}\right), \\
N(\mu)_{i}^{\exp }=N_{i}^{\mathrm{non}-\mathrm{QE}} \cdot\left(1+\sum_{j=4,6,7} f(\mu)_{j}^{i} \cdot \epsilon_{j}\right)+N_{i}^{\mathrm{QE}} \cdot\left(1+\sum_{j=4,5,7} f(\mu)_{j}^{i} \cdot \epsilon_{j}\right) .
\end{gathered}
$$

$f_{j}^{i}$ : fractional change in the predicted event rate in the i-th bin due to the variation of the parameter $\epsilon_{j}$

Systematic Errors
5 \% BG (Overall)
$\epsilon_{j}$ : systematic error parameters varied freely to minimize $\chi^{2}$
5 \% BG (Energy Dep.)
5 \% Signal Efficiency
20 \% Separation of QE/nQE

## Example of the case where the octant degeneracy is resolved

HK(0.54Mt) only at Kamioka

$2 \mathrm{HKs}(0.27 \times 2 \mathrm{Mt})$ at Kamioka and Korea
Kamioka 0.27Mton + Korea 0.27Mton detectors, $v 4 y r+\bar{v} 4 y r 4 M W$ beams


Octant degeneracy is resolved!

$$
\sin ^{2} \theta_{23}=0.60 \text { (true) }
$$

(Despite that sign $\Delta \mathrm{m}^{2}$ degeneracy is not completely resolved $\rightarrow$ decoupling of degeneracy)

Parameter regions where the octant degeneracy can be resolved


Dependence on the mass hierarchy is weak

## Comparing two methods...

T2K II + Reactor


T2KK


$\sin ^{2} 2 \theta_{13} \geq 0.04-0.06:$ T2KII + Reactor is better $\sin ^{2} 2 \theta_{13} \lesssim 0.04-0.06:$ T2KK is better

## Summary

- Octant degeneracy exist if $\theta_{23}$ is different from $\pi / 4$
- MethodI:Superbeam (T2K II) + Reactor can resolve the octant degeneracy for $\sin ^{2} 2 \theta_{23} \approx 0.96$ (0.99) if $\sin ^{2} 2 \theta_{13} \approx 0.05$ (0.1) but not work for small $\theta_{13}$
- Method II: Superbeam with 2 Far Detector System (T2KK) can Resolve the octant degeneracy for $\sin ^{2} 2 \theta_{23}$ $<0.97$ even for very small $\theta_{13}$
- Both Methods are Complementary: For larger (smaller) $\theta_{13}, \sin ^{2} 2 \theta_{13} \gtrsim(\Sigma) 0.05$, Method I (II) would be better

