# QCD Simulations at Realistic Quark Masses: Probing the Chiral Limit

G. Schierholz

Deutsches Elektronen-Synchrotron DESY

– QCDSF Collaboration –



Special mention:

M. Göckeler, T. Hemmert, R. Horsley, Y. Nakamura, D. Pleiter, P.E.L. Rakow, W. Schroers, T. Streuer, H. Stüben and J. Zanotti

#### Objective

Solve QCD and probe the limits of the Standard Model · · ·

• Parameters of QCD

- $\Lambda_{QCD}$  resp.  $lpha_s(Q^2)$
- Quark masses
- $\theta$  angle

- QCD in the wider world
- How does QCD work?
- Fundamental properties

- CKM matrix
- Hadron structure
- Spectroscopy
- $\chi SB$
- Confinement

 $\cdots$  in concert with Exp & Phen

#### Problem: Chiral Extrapolation



Recently

Need to reduce (scale) error to a few %

Outline

**Lattice Simulations** 

**Pion Sector** 

**Nucleon Sector** 

**Miscellaneous** 

**Conclusions & Outlook** 

Lattice Simulations

Action

$$N_f = 2$$

 $S = S_G + S_F$ 

$$S_G = \beta \sum_{x,\mu < \nu} \left( 1 - \frac{1}{3} \operatorname{Re} \operatorname{Tr} U_{\mu\nu}(x) \right)$$

$$S_{F} = \sum_{x} \left\{ \bar{\psi}(x)\psi(x) - \kappa \,\bar{\psi}(x)U_{\mu}^{\dagger}(x-\hat{\mu})[1+\gamma_{\mu}]\psi(x-\hat{\mu}) - \kappa \,\bar{\psi}(x)U_{\mu}(x)[1-\gamma_{\mu}]\psi(x+\hat{\mu}) - \frac{1}{2}\kappa \,c_{SW} \,g \,\bar{\psi}(x)\sigma_{\mu\nu}F_{\mu\nu}(x)\psi(x) \right\}$$

$$\Im \ \partial_\mu A_\mu^{
m imp} = 2 m_q P$$

# **Clover Fermions**

#### Advantages

- Local
- Transfer matrix
- O(a) improved
- Flavor symmetry

Prerequisite to making contact with  $SU(2)\ {\rm ChPT}$ 

- Finite size corrections
- Chiral extrapolation
- Determination of low-energy constants

• Fast to simulate

#### Cost of Simulation

1000 Configurations



 $\propto L^{4.8}\,(m_\pi/m_
ho)^{-3.6}\,(r_0/a)^{0.9}$ 

Hasenbusch, QCDSF, Lüscher, Urbach et al., · · ·

# Compared to $\cdots$



Clark

#### Parameters



For gauge field sampling we use 'ordinary' HMC algorithm with Hasenbusch integration + 3 time scales

# Obstructions ?







Aoki phase



# Landscape



Minimal pion mass : 
$$m_{\pi}(L) = \frac{3}{2f_0^2 L^3} \left(1 + \frac{2}{4\pi f_0^2 L^2} 2.837\right)^{-1}$$

Leutwyler Hasenfratz & Niedermayer

#### Effect of Unquenching ?

Vector Ward Identity ?



Dürr

**Pion Sector** 





NLO

$$m_{PS}^{2} = m_{0}^{2} \left[ 1 + \frac{1}{2} x \,\hat{l}_{3} + O(x^{2}) \right]$$
$$\frac{m_{PS} - m_{PS}(L)}{m_{PS}} = -\sum_{|\vec{n}|\neq 0} \frac{x}{2\lambda} \left[ I_{m_{PS}}^{(2)}(\lambda) + x I_{m_{PS}}^{(4)}(\lambda) \right]$$



$$m_0^2 = 2\Sigma m_q, \ x = \frac{m_0^2}{16\pi^2 f_0^2}, \ \lambda = m_{PS} |\vec{n}| L$$
$$\hat{l}_i = \ln \frac{\Lambda_i^2}{m_0^2}$$

No  $1^{st}$  order phase transition or Aoki phase !

$$\begin{split} I_{m_{PS}}^{(2)}(x) &= -B^0(x) \\ I_{m_{PS}}^{(4)}(x) &= \left( -\frac{55}{18} + 4\bar{l}_1 + \frac{8}{3}\bar{l}_2 - \frac{5}{2}\bar{l}_3 - 2\bar{l}_4 \right) B^0(x) \\ &+ \left( \frac{112}{9} - \frac{8}{3}\bar{l}_1 - \frac{32}{3}\bar{l}_2 \right) B^2(x) + S_{m_{PS}}^{(4)}(x) \end{split}$$

$$S_{m_{PS}}^{(4)}(x) = \frac{13}{3}g_0 B^0(x) - \frac{1}{3}\left(40g_0 + 32g_1 + 26g_2\right) B^2 + \cdots$$

$$B^0(x) = 2K_1(x)\,, \quad B^2(x) = 2K_2(x)/x\,, \quad ar{l}_i = \ln rac{\Lambda_i^2}{m_{PS}^2}$$

 $\Lambda_i, \ g_i \ {\sf from \ hep-lat}/05030142$ 



$$r_0 f_0 = 0.179(2), \quad r_0 \Lambda_3 = 1.82(7)$$

#### Pion Decay Constant



$$\begin{split} I_{f_{PS}}^{(2)}(x) &= -2B^0(x) \\ I_{f_{PS}}^{(4)}(x) &= \left(-\frac{7}{9} + 2\bar{l}_1 + \frac{4}{3}\bar{l}_2 - 3\bar{l}_4\right)B^0(x) \\ &+ \left(\frac{112}{9} - \frac{8}{3}\bar{l}_1 - \frac{32}{3}\bar{l}_2\right)B^2(x) + S_{f_{PS}}^{(4)}(x) \end{split}$$

$$S_{f_{PS}}^{(4)}(x) = \frac{1}{6} \left( 8g_0 - 13g_1 \right) B^0(x) - \frac{1}{3} \left( 40g_0 - 12g_1 - 8g_2 - 13g_3 \right) B^2 + \cdots$$

Colangelo, Dürr & Haefeli

#### Partially Quenched

$$m_{PS} \equiv m_{PS}^{SS} \to m_{PS}^{AB}, \quad f_{PS} \equiv f_{PS}^{SS} \to f_{PS}^{AB}, \qquad A, B \in \{V, S | V \neq S\}$$



Sharpe

**Nucleon Sector** 

### Nucleon Mass



$$m_{N} = m_{0} - 4c_{1}m_{PS}^{2} - \frac{3g_{A}^{0\,2}}{32\pi f_{0}^{2}}m_{PS}^{3} + \left[e_{1}(\mu) - \frac{3}{64\pi^{2}f_{0}^{2}}\left(\frac{g_{A}^{0\,2}}{m_{0}} - \frac{c_{2}}{2}\right) - \frac{3g_{A}^{0\,2}}{32\pi^{2}f_{0}^{2}}\left(\frac{g_{A}^{0\,2}}{m_{0}} - 8c_{1} + c_{2} + 4c_{3}\right)\ln\frac{m_{PS}}{\mu}\right]m_{PS}^{4} + \frac{3g_{A}^{0\,2}}{256\pi f_{0}^{2}m_{0}^{2}}m_{PS}^{5} + O(m_{PS}^{6})$$

$$\begin{split} m_N - m_N(L) &= -\frac{3g_A^{0\,2}m_0m_{PS}^2}{16\pi^2 f_0^2} \sum_{|\vec{n}|\neq 0} \int_0^\infty dz K_0 \left(\sqrt{m_0^2 z^2 + m_{PS}^2(1-z)} |\vec{n}|L\right) \\ &- \frac{3m_{PS}^4}{4\pi^2 f_0^2} \sum_{|\vec{n}|\neq 0} \left[ (2c_1 - c_3) \frac{K_1(m_{PS}|\vec{n}|L)}{m_{PS}|\vec{n}|L} + c_2 \frac{K_2(m_{PS}|\vec{n}|L)}{(m_{PS}|\vec{n}|L)^2} \right] + O(m_{PS}^5) \end{split}$$

# Axial Coupling





 $\chi {\sf PT}~O(p^3)$ 

 $68.3\%~{\rm CL}$ 

Miscellaneous

# Rho Mass



Not FS corrected

# Delta Mass



Not FS corrected

**Conclusions & Outlook** 

- Simulations at pion masses of O(300) MeV with Wilson-type fermions feasible now
- Extrapolation to chiral limit and infinite volume greatly improved

- Improvement of algorithms
- Increase of computing power

FS corrections surprisingly well described by ChPT

• First meaningful lattice determination of low energy constants : Preliminary !

| $r_0$      | $f_0$     | $\Lambda_3$ | $\Lambda_4$  |
|------------|-----------|-------------|--------------|
| 0.45(3) fm | 79(5) MeV | 0.80(5) GeV | 1.46(10) GeV |

• Major investment in FS corrections (including partially quenched data) and  $\delta$  expansion needed